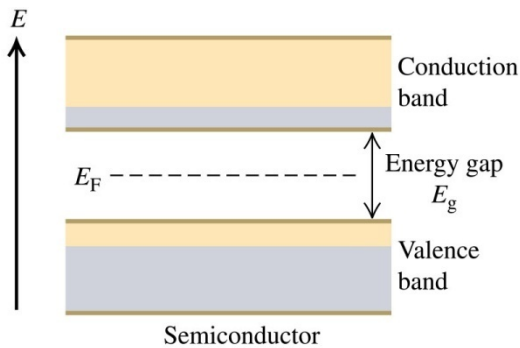




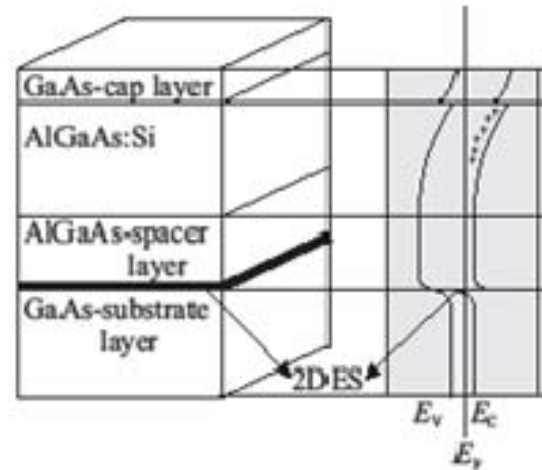
COULOMB EFFECTS IN ULTRA-SMALL ELECTRONIC DEVICES

ANDREI MANOLESCU, ASSOCIATE PROFESSOR
SCHOOL OF SCIENCE AND ENGINEERING | RU LECTURE MARATHON

Nanoelectronic devices



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GaAs/AlGaAs heterostructure

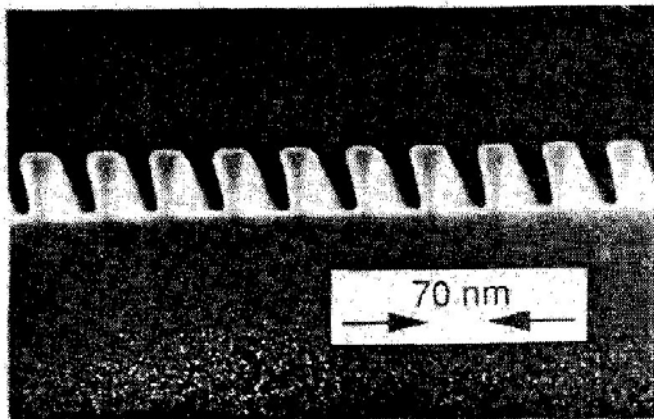
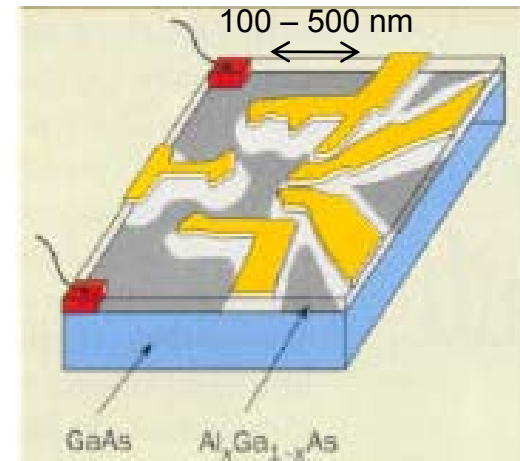


Fig. 2. Cross-sectional view of the reverse-mesa etched QWJs taken by a high resolution SEM.

Quantum wires



Quantum dots

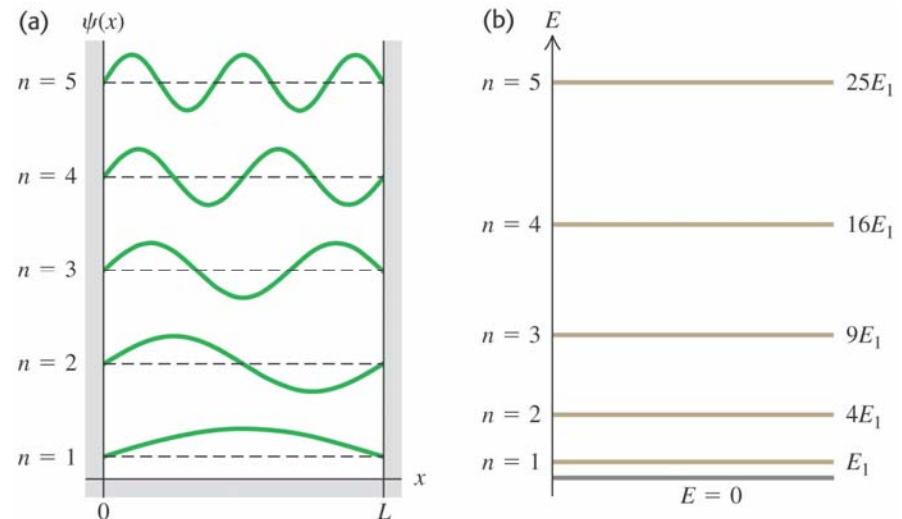
Quantum mechanics

Wave functions: $\Psi(\vec{r}, t)$

Schrödinger equation: $i\hbar \frac{\partial}{\partial t} \Psi(\vec{r}, t) = H(t) \Psi(\vec{r}, t)$

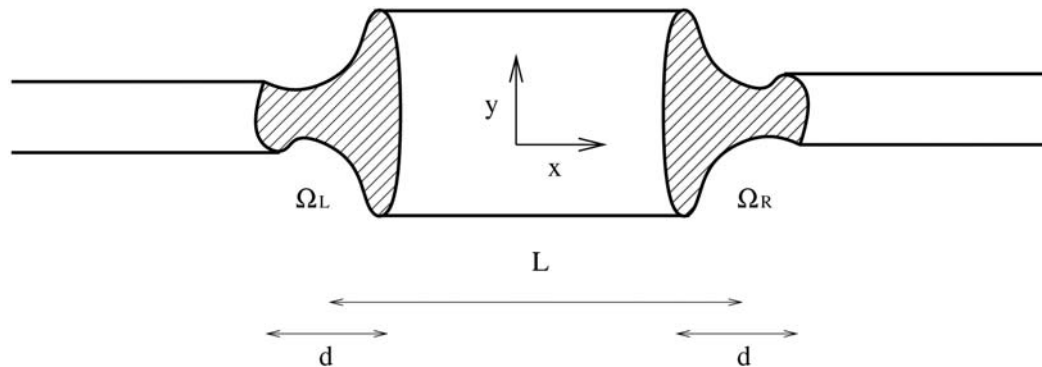
Hamiltonian (energy) operator: $H(t) = -\frac{\hbar^2}{2m} \Delta + U(\vec{r}, t)$

Example: The quantum well



$$\psi_n(x) = \sqrt{\frac{2}{L}} \sin \frac{n\pi x}{L} \quad (n = 1, 2, 3, \dots)$$

Open quantum dots in contact with leads



Unknown wave functions of the electrons in the sample

Statistical description: wave functions ψ_n with probabilities p_n

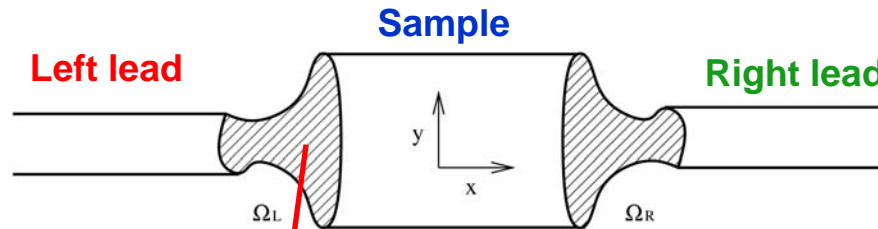
Statistical operator $\rho = \sum_n p_n | \psi_n \rangle \langle \psi_n |$

Expected values of physical observables:

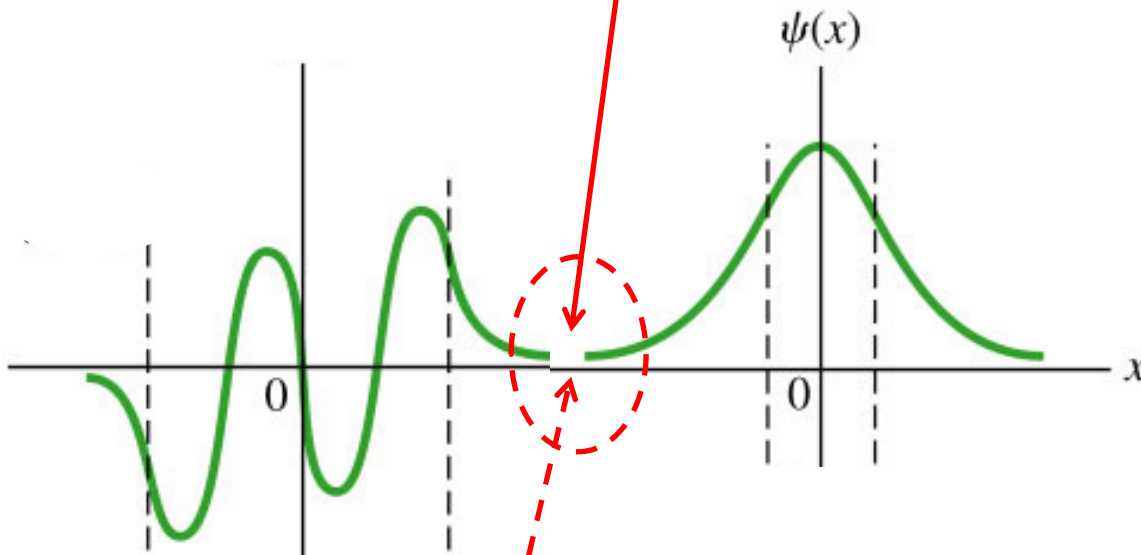
$$\langle A \rangle_n = \langle \psi_n | A | \psi_n \rangle \rightarrow \langle A \rangle = \sum_n p_n \langle \psi_n | A | \psi_n \rangle$$

$$\Rightarrow \langle A \rangle = \text{Tr}(A\rho) = \sum_n (A\rho)_{nn}$$

The leads-sample contacts

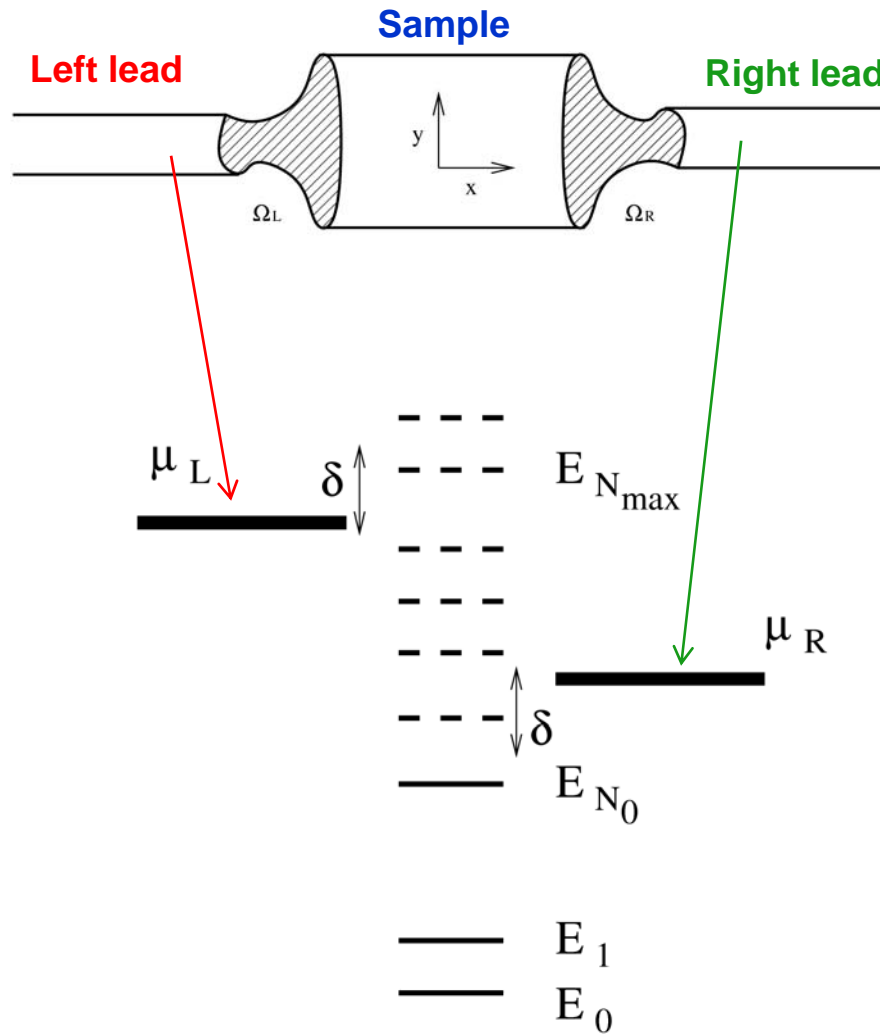


Overlap of wave functions



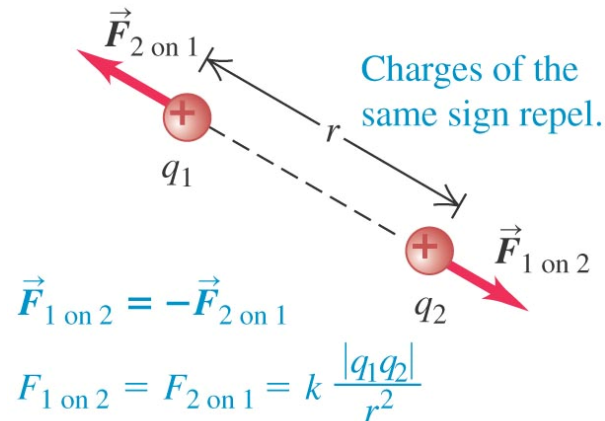
$$T_{qnl} \propto \int_{\Omega_l} d\vec{r} \int_{\Omega_l} d\vec{r}' \psi_{ql}^*(\vec{r}) g(\vec{r}, \vec{r}') \varphi_n(\vec{r}') \rightarrow \text{Generic form}$$

The external field (bias)



$\mu_{L,R}$ chemical potentials

Coulomb interaction and many-body effects



$$N \text{ electrons in the sample} \Rightarrow H_{\text{Sample}} = \sum_n T_n + \frac{1}{2} \sum_{n \neq m} \frac{e^2}{|\vec{r}_n - \vec{r}_m|}$$

$$V_{nm,pq} = \int d\vec{r} d\vec{r}' \psi_n^*(\vec{r}) \psi_m^*(\vec{r}') \frac{e^2}{|\vec{r}_n - \vec{r}_m|} \psi_p(\vec{r}) \psi_q(\vec{r}')$$

The equation for the statistical operator (generalized master equation)

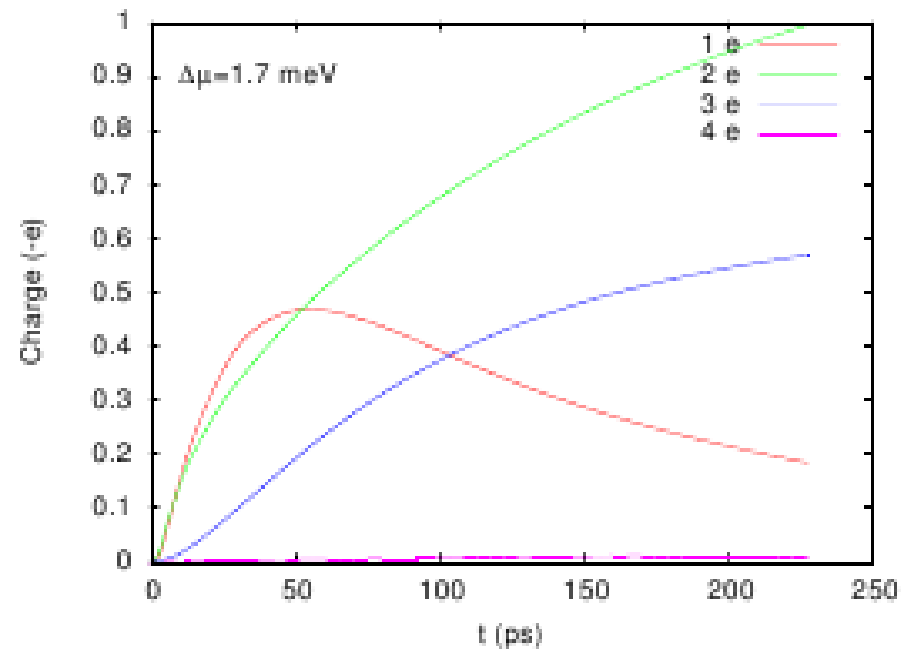
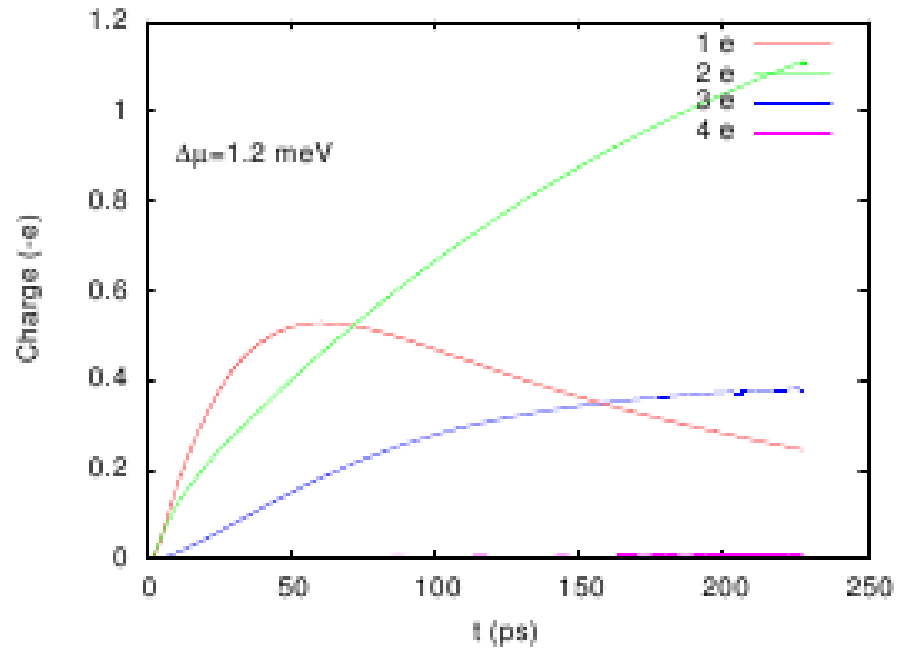
$$\dot{\rho}_S(t) = -\frac{i}{\hbar}[H_S, \rho_S] - \frac{1}{\hbar^2} \sum_l \int dq \chi_l(t) \left\{ [T_{ql}, \Omega_{ql}(t)] + h.c. \right\}$$

$$\Omega_{ql}(t) = e^{-itH_S/\hbar} \Pi_{ql}(t) e^{itH_S/\hbar} e^{-it\varepsilon_l(q)/\hbar}$$

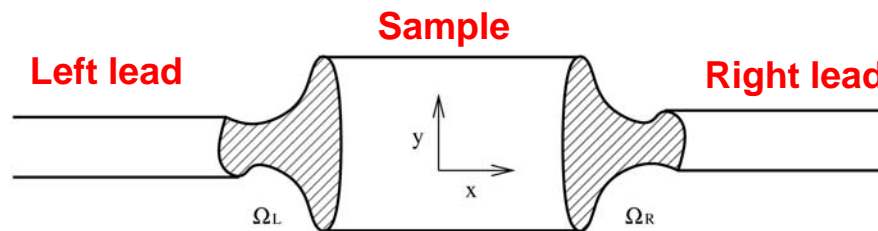
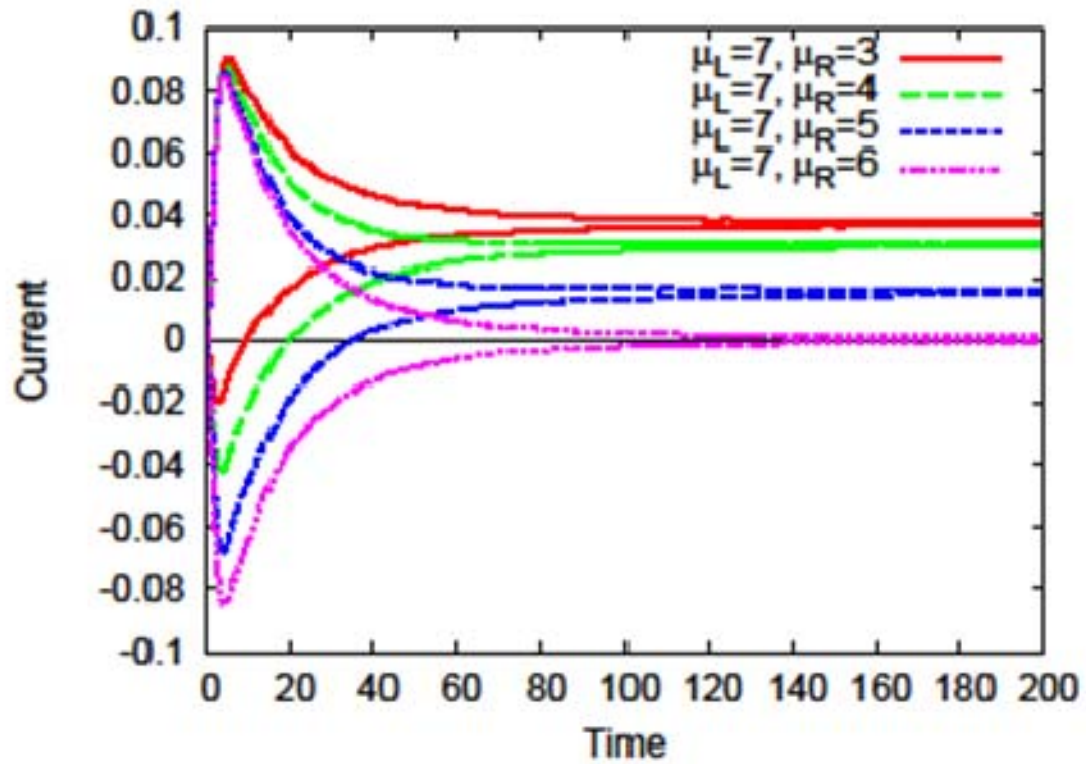
$$\Pi_{ql}(t) = \int_0^t ds \chi_l(s) e^{itH_S/\hbar} \times$$

$$\times \left\{ T_{ql}^\dagger \rho_S(s) [1 - f(\varepsilon_l(q))] - \rho_S(s) T_{ql}^\dagger f(\varepsilon_l(q)) \right\} e^{-isH_S/\hbar} e^{is\varepsilon_l(q)/\hbar}$$

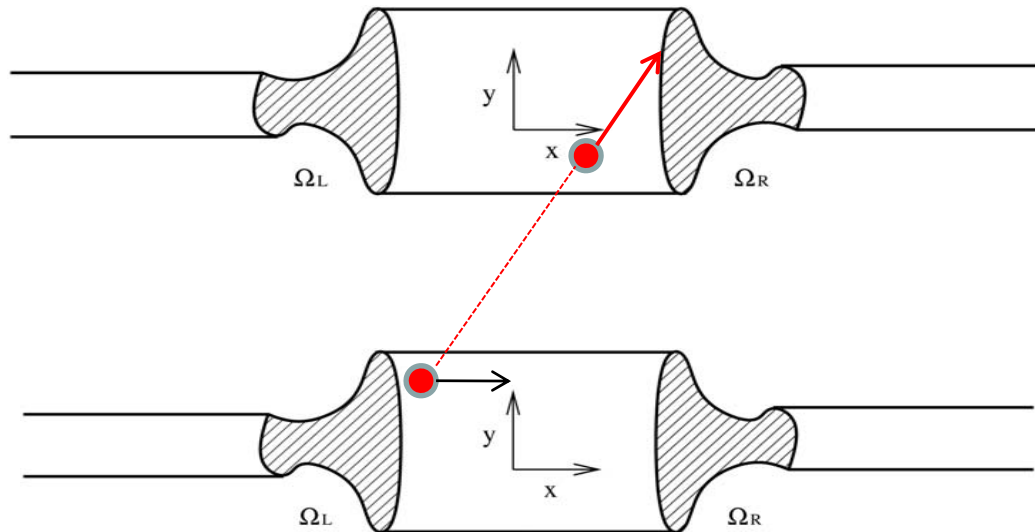
Charging of the sample



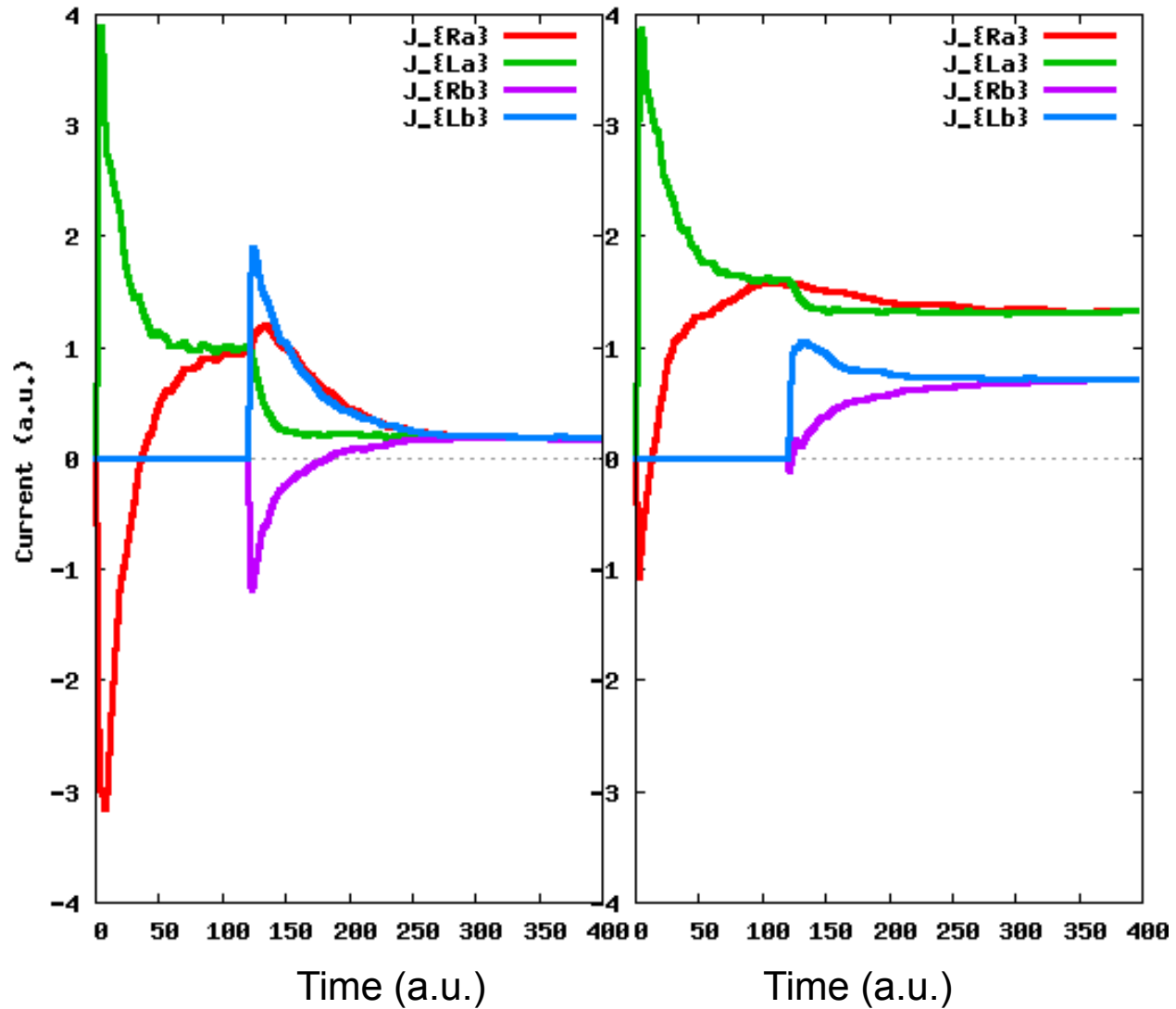
Currents



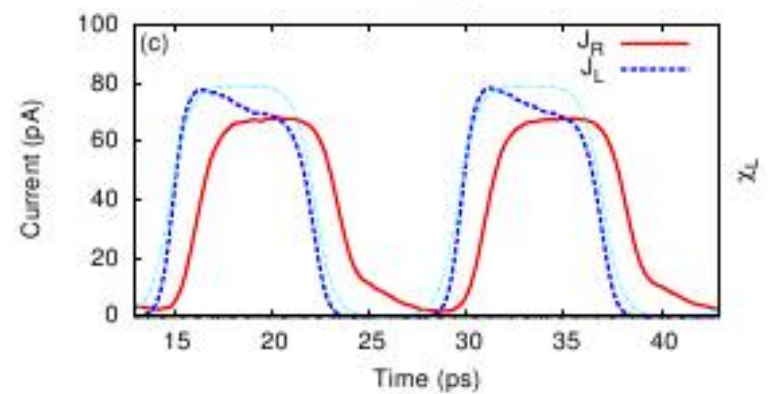
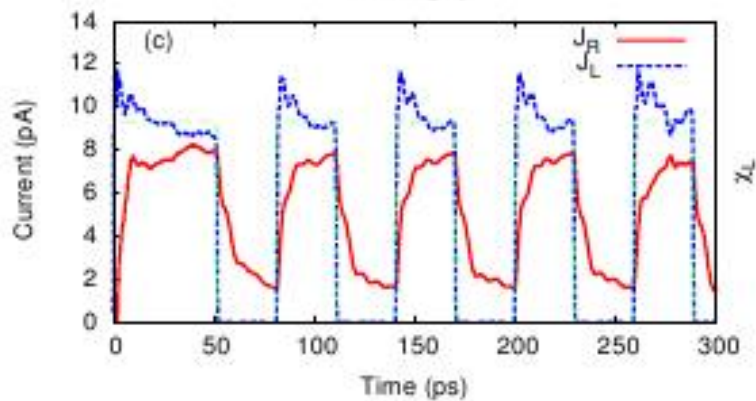
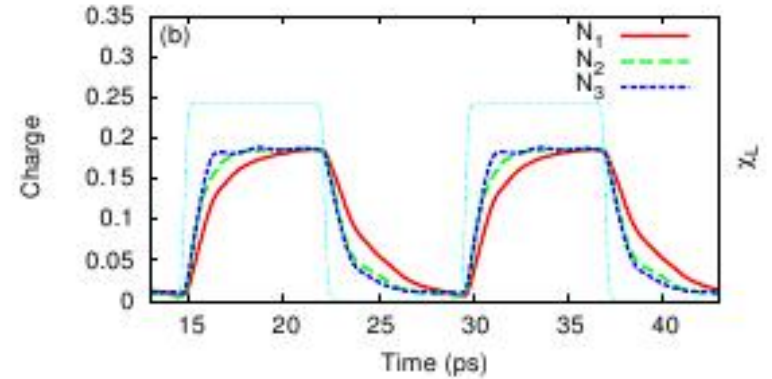
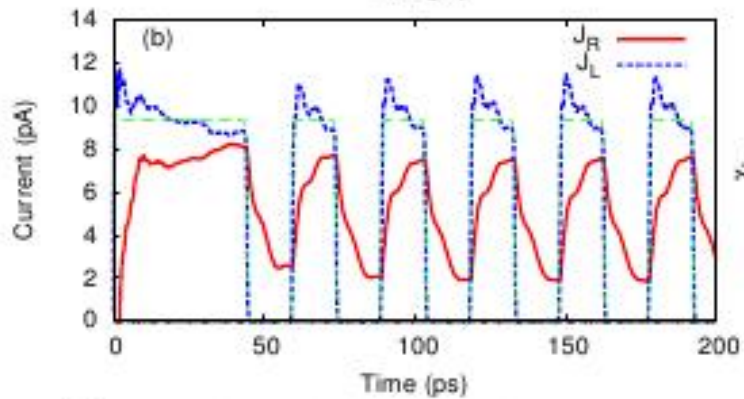
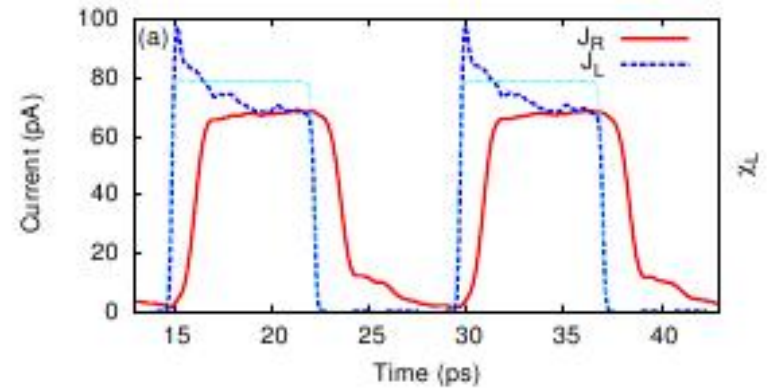
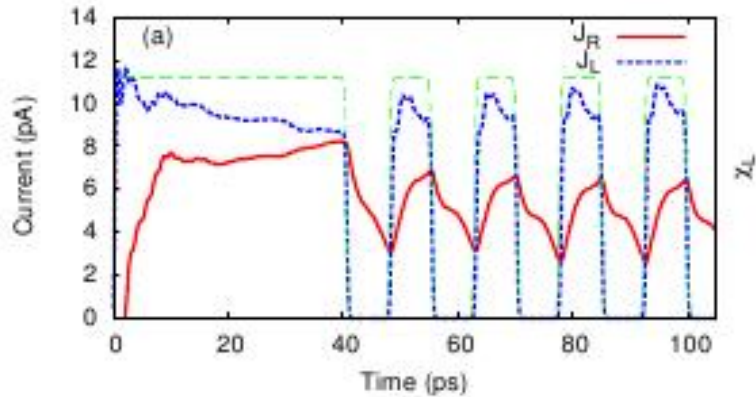
Double system



Induced currents



Electric pulses



The team



Viðar Guðmundsson,
Professor HÍ



Valeriu Moldoveanu,
Scientist, Romania



Cosmin Gainer,
PhD student HÍ



Kristinn Torfason,
PhD student HR



Andrei, speaker